



General relativity, astronomy and the universe:
the first hundred years

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Waterford Institute of Technology

Irish National Astronomy Meeting UCD 7-9/9/16

Overview



Einstein in Berlin (1916)

Relativity

The special theory of relativity

The general theory of relativity

Three classic astronomical tests

The perihelion of Mercury; the bending of starlight

The gravitational redshift

Three cosmological predictions

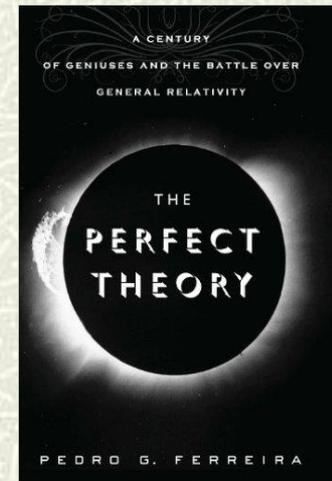
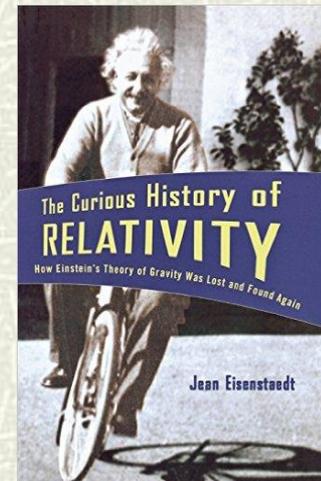
Black holes; gravitational waves

The dynamic universe

The curious history of general relativity

Low watermark period (1930-60)

Astronomy catches up - the golden period



The special theory of relativity (1905)



Two principles

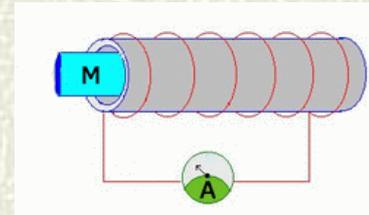
Laws of all physics identical for observers in relative uniform motion

Speed of light in vacuum identical for observers in relative uniform motion

Implications

Intervals in distance and time not universal

Experienced differently by bodies in relative uniform motion



Predictions (high-speed bodies)

Length contraction : time dilation

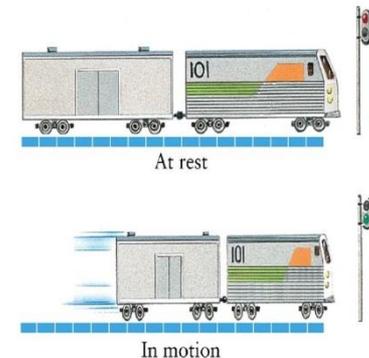
Mass increase; mass-energy equivalence

Space-time (Minkowski 1908)

Space-time invariant for observers in inertial frames

$$ds^2 = dx^2 + dy^2 + dz^2 - c^2 dt^2$$

Special Relativity: Length Contraction



The general theory of relativity (1916)

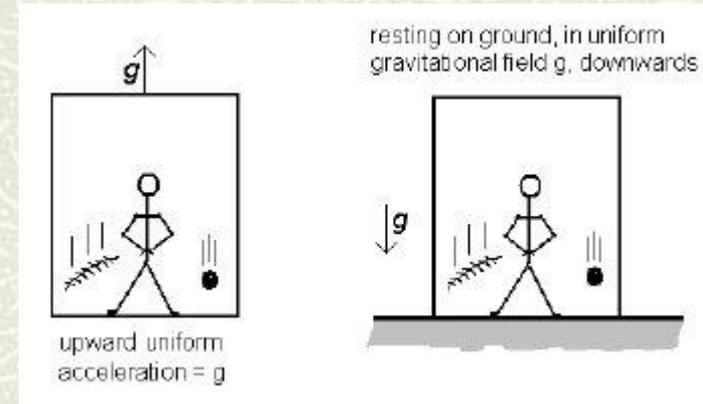
Extending the special theory (1907-)

Relativity and gravity?

Relativity and accelerated motion?

The principle of equivalence

Equivalence of gravity and acceleration



Mach's principle

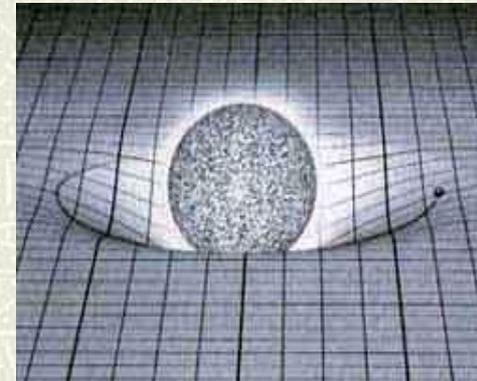
Inertial mass defined relative to matter

"No such thing as empty space"

A long road (1907-1915)

Curvilinear geometry, tensor algebra

Gravity = curvature of space-time



The field equations of GR (1915)



$$G_{\mu\nu} = - \frac{8\pi G}{c^4} T_{\mu\nu}$$

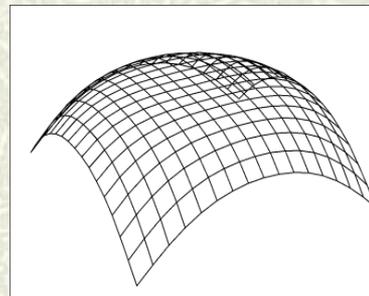


10 non-linear differential equations relating the spacetime metric to the density and flow of energy and momentum

SR $ds^2 = dx^2 + dy^2 + dz^2 - c^2 dt^2$

$$ds^2 = \sum_{\mu, \nu=1}^4 n_{\mu\nu} dx^\mu dx^\nu$$

$$n_{\mu\nu} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$



GR $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$

$$ds^2 = \sum_{\mu, \nu=1}^4 g_{\mu\nu} dx^\mu dx^\nu$$

$g_{\mu\nu}$: variables determined by matter

Three astronomical tests (Einstein, 1916)

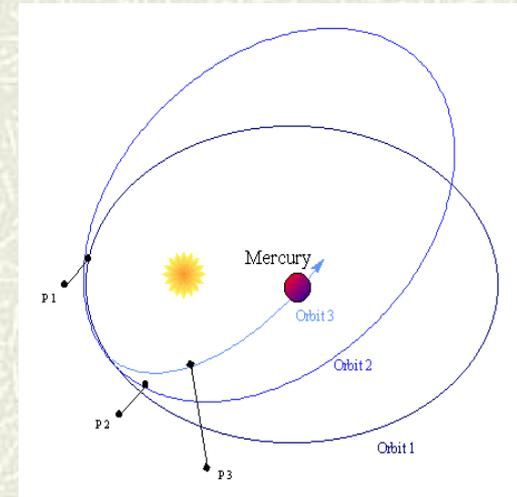
Small deviations from Newton's gravity

Different in principle

The perihelion of Mercury

Well-known anomaly in Mercury's orbit (43" per century)

Postdicted by GR (1916)



The bending of starlight by the sun (1.7")

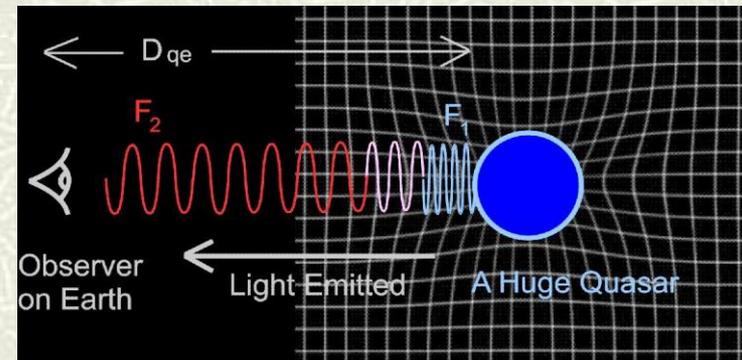
Eclipse expeditions of Eddington and Dyson (1919)

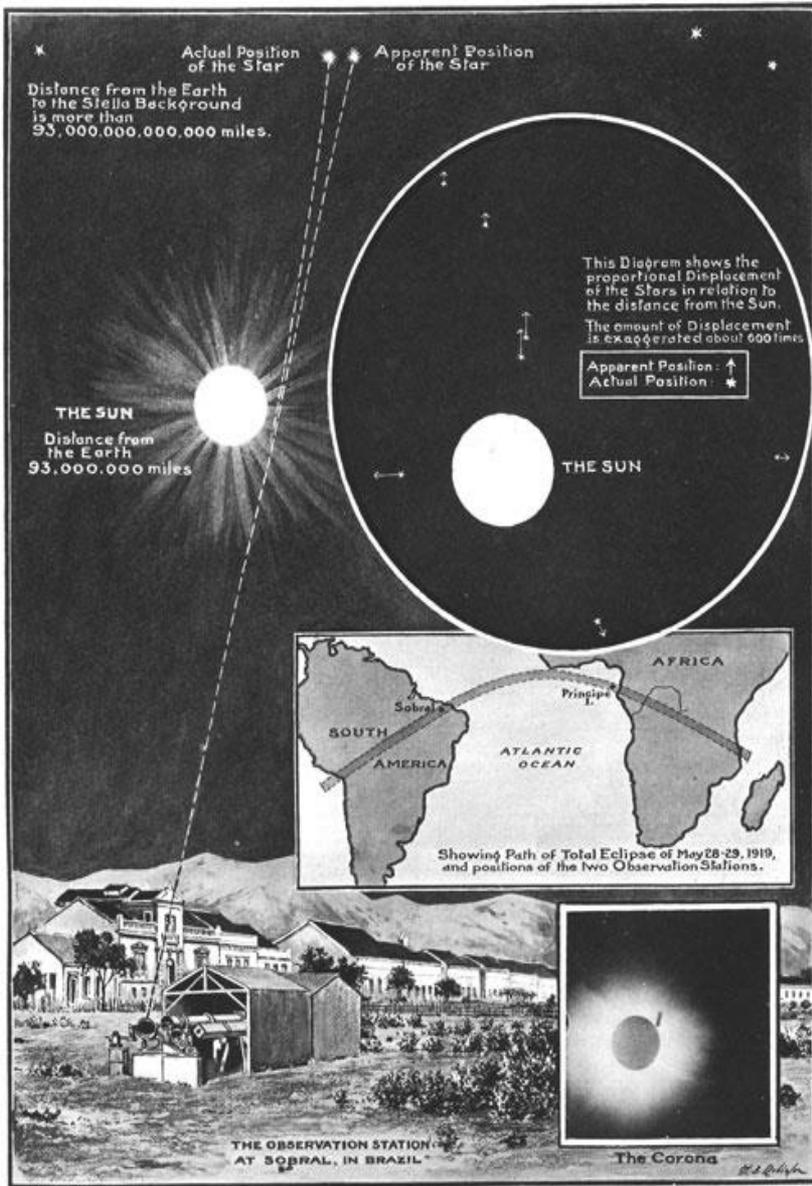
Successful measurement (large error margin)

Gravitational redshift

Time dilation in strong gravitational field

Light from a massive star redshifted?





Eclipse Results (1919)

Sobral: 1.98" +/- 0.16

Principe: 1.7" +/- 0.4

Einstein famous (1919)

LIGHTS ALL ASKEW IN THE HEAVENS

Men of Science More or Less
Agog Over Results of Eclipse
Observations.

EINSTEIN THEORY TRIUMPHS

Stars Not Where They Seemed
or Were Calculated to be,
but Nobody Need Worry.

A BOOK FOR 12 WISE MEN

No More in All the World Could
Comprehend It, Said Einstein When
His Daring Publishers Accepted It.

Asymmetric controversy (Collins and Pinch 1970s)

Claim of bias; rebutted by astronomers (RAS)

Einstein's reaction

Albert Einstein, The Times (Nov 28th 1919)

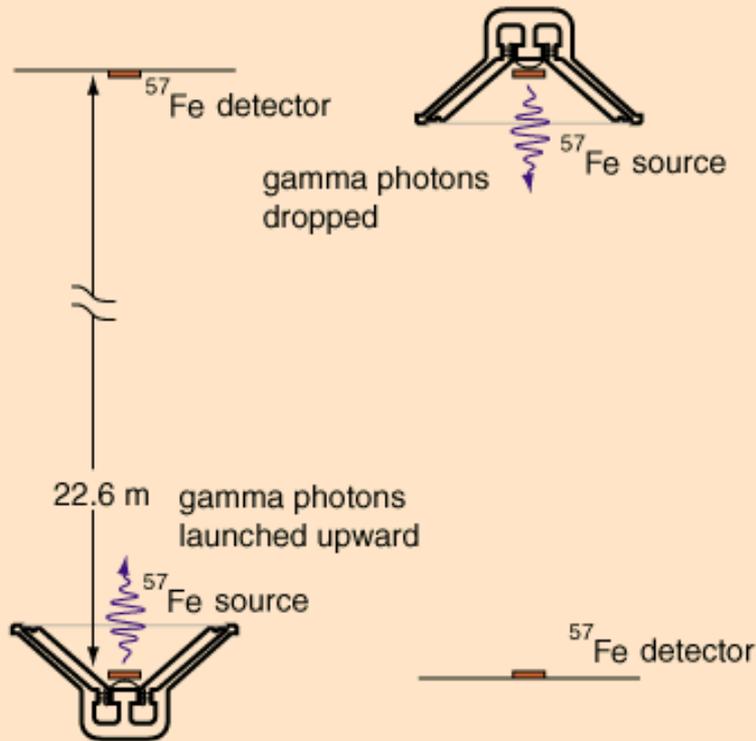
The new theory of gravitation diverges considerably, as regards principles, from Newton's theory. But its practical results agree so nearly with those of Newton's theory that it is difficult to find criteria for distinguishing them which are accessible to experience. Such have been discovered so far:

1. In the revolution of the ellipses of the planetary orbits round the sun (confirmed in the case of Mercury).
2. In the curving of light rays by the action of gravitational fields (confirmed by the English photographs of eclipses).
3. In a displacement of the spectral lines toward the red end of the spectrum in the case of light transmitted to us from stars of considerable magnitude (unconfirmed so far).*

Let no one suppose, however, that the mighty work of Newton can really be superseded by this or any other theory. His great and lucid ideas will retain their unique significance for all time as the foundation of our whole modern conceptual structure in the sphere of natural philosophy.

3. Gravitational redshift (1959)

Harvard Tower Experiment



In just 22.6 meters, the fractional gravitational red shift given by

$$\nu = \nu_0 \left[1 + \frac{gh}{c^2} \right]$$

is just 4.92×10^{-15} , but the Mössbauer effect with the 14.4 keV gamma ray from iron-57 has a high enough resolution to detect that difference. In the early 60's physicists Pound, Rebka, and Snyder at the Jefferson Physical Laboratory at Harvard measured the shift to within 1% of the predicted shift.

- Sirius B (Adams, 1925)
- Gravity Probe A (1970s)

Cosmic prediction I: Black Holes

Schwarzschild (1916)

Exact solution for the field equations

Body of spherical symmetry

Enigma

Solution becomes singular at $r = 2GM/c^2$

Space closed up around mass?

Rejected

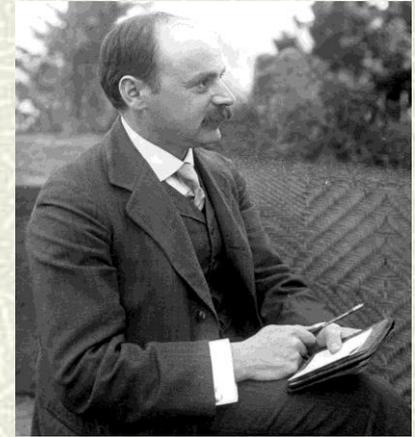
Co-ordinate problem (Eddington)

Prevented by internal pressure (Einstein 1922)

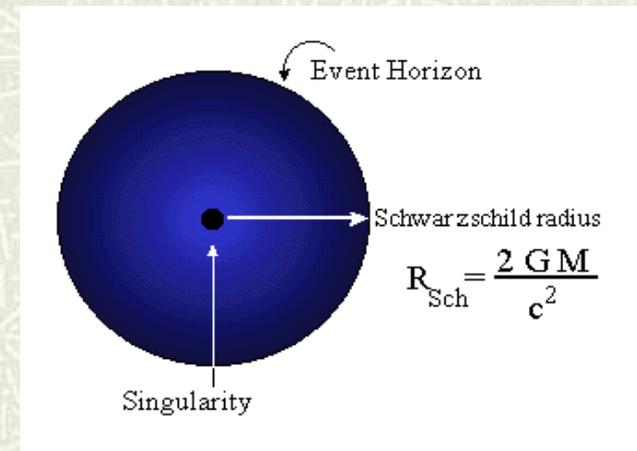
Physical reality?

Collapse of sun? Anderson (UCG)

Collapse of large stellar ensemble : Lodge (Liverpool)



Karl Schwarzschild (1873–1916)



The physics of black holes

Chandrasekhar (1931)

The physics of white dwarf stars (quantum degeneracy)

SR: collapse to infinite density for $M > 1.4 M_{\odot}$

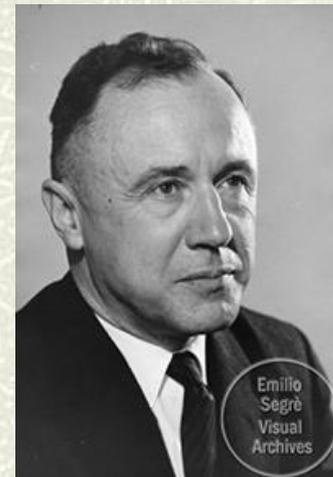
Rejected by Eddington, community



Oppenheimer (1939,40)

GR: Continued stellar collapse for $M > 3 M_{\odot}$

Rejected by Einstein (1939)



Wheeler, Thorne, Zeldovitch (1960s)

Numerical solutions of the field equations

Simulation of stellar collapse

Penrose (1965)

No avoiding BH singularity

Black Holes: Observation

Compact astronomical objects (1960s)

Quasars: small, distant sources of incredible energy (1963)

Pulsars: rapidly rotating neutron stars (1967)

X-ray binaries

Cygnus X-1 (1964)

Matter pulled from star into massive companion emits X-rays

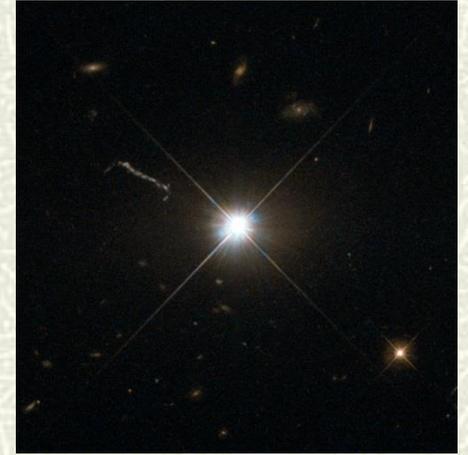
Orbit studies

Supermassive BH at centre of MW? (1990s)

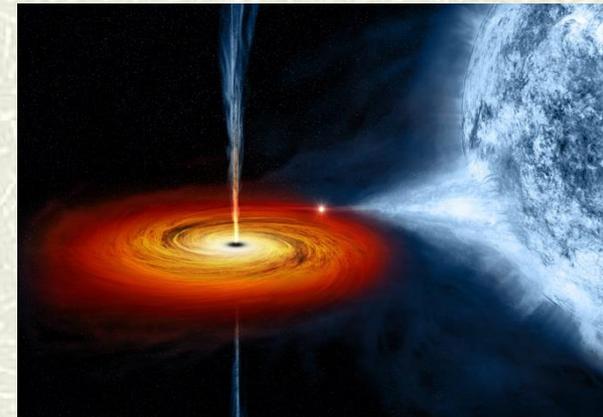
Supermassive BH at centre of many galaxies (2000-)

2015-16

Gravitational waves from binary BH system!



Quasar 3C273



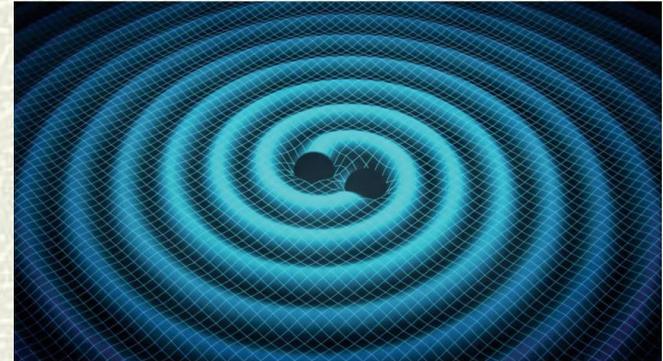
Cygnus X-1 (1964)

Cosmic prediction II: gravitational waves

Einstein (1916, 18)

Linearized wave-like solutions of GFE

Cosmic events cause tiny ripples in spacetime?



Einstein and Rosen (1936, 37)

Cylindrical wave solutions - no gravitational waves (1936)

Corrected with assistance from HP Robertson (1937)

Wheeler (1960s)

Numerical wave solutions

Weber bars (1960s)

Reports signal of gravitational waves

Not reproduced, not accepted by community

Joseph Weber



Gravitational Waves: Observation

Binary pulsar PSR1913+16

Hulse-Taylor (1974)

Decrease in orbital period exactly as predicted

Direct measurement?

Interferometers: 1980-

Interferometers with 4 km arms (LIGO, VIRGO)

Advanced LIGO (2015)

Clear signal (September 2015)

Double whammy

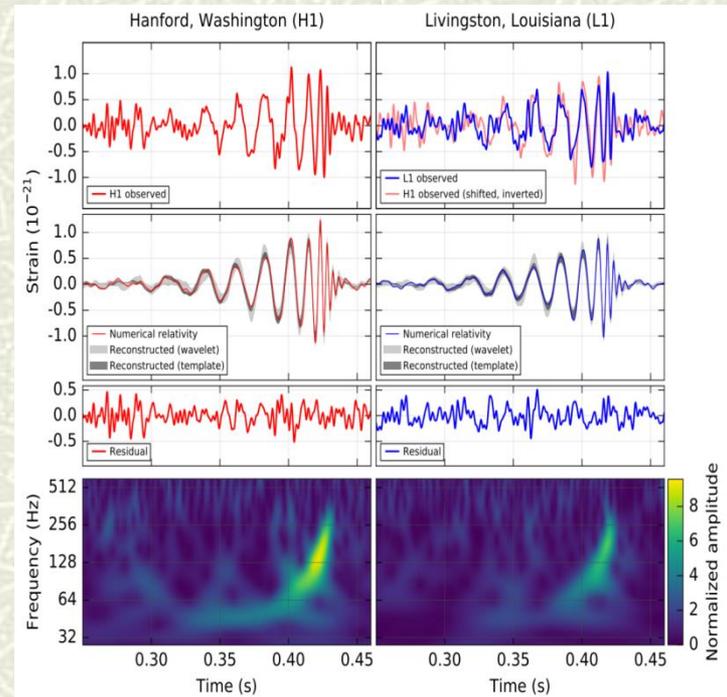
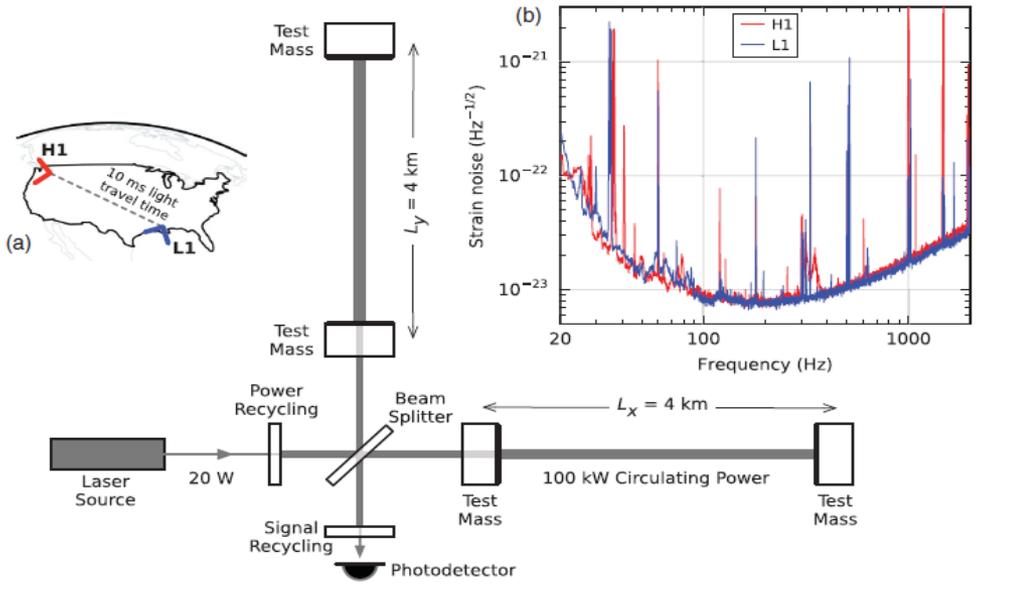
Exact match with merging BHs

$29 M_{\odot}$, $36 M_{\odot}$; 1.3 billion LY away



Hulse-Taylor pulsar





1. Shape of waveform
2. Frequency of orbit

BBH !

061102 (2016)

Selected for a Viewpoint in *Physics*
 PHYSICAL REVIEW LETTERS

week
 12 FEBRU



Observation of Gravitational Waves from a Binary Black Hole Merger

B. P. Abbott *et al.**

(LIGO Scientific Collaboration and Virgo Collaboration)

(Received 21 January 2016; published 11 February 2016)

On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with a peak gravitational-wave strain of 1.0×10^{-21} . It matches the waveform predicted by general relativity for the inspiral and merger of a pair of black holes and the ringdown of the resulting single black hole. The signal was observed with a matched-filter signal-to-noise ratio of 24 and a false alarm rate estimated to be less than 1 event per 203 000 years, equivalent to a significance greater than 5.1σ . The source lies at a luminosity distance of 410^{+160}_{-180} Mpc corresponding to a redshift $z = 0.09^{+0.03}_{-0.04}$. In the source frame, the initial black hole masses are $36^{+4}_{-4} M_{\odot}$ and $29^{+4}_{-4} M_{\odot}$, and the final black hole mass is $62^{+4}_{-4} M_{\odot}$, with $3.0^{+0.5}_{-0.5} M_{\odot} c^2$ radiated in gravitational waves. All uncertainties define 90% credible intervals. These observations demonstrate the existence of binary stellar-mass black hole systems. This is the first direct detection of gravitational waves and the first observation of a binary black hole merger.

Cosmic prediction III: The expanding universe

Einstein: apply GR to the Universe (1917)

Ultimate test for new theory of gravitation

Assumptions

Static, uniform distribution of matter

Mach's principle: metric tensor to vanish at infinity

Boundary problem

Assume cosmos of closed curvature

But...no consistent solution

New term in field equations!

Cosmic constant - anti-gravity term

Radius and density defined by λ

$$G_{\mu\nu} = -\kappa T_{\mu\nu}$$

$$G_{\mu\nu} + \lambda g_{\mu\nu} = -\kappa T_{\mu\nu}$$

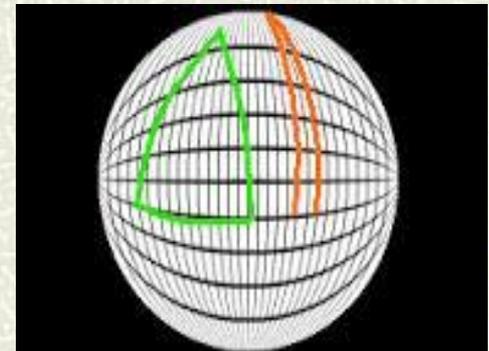
Doc. 43

Cosmological Considerations in the General Theory of Relativity

This translation by W. Perrett and G. B. Jeffery is reprinted from H. A. Lorentz et al., *The Principle of Relativity* (Dover, 1952), pp. 175–188.

It is well known that Poisson's equation

$$\nabla^2\phi = 4\pi K\rho \quad (1)$$
 in combination with the equations of motion of a material point is not as yet a perfect substitute for Newton's theory of action at a distance. There is still to be taken into account the condition that at spatial infinity the potential ϕ tends



$$\lambda = \frac{\kappa\rho}{2} = \frac{1}{R^2}$$

The dynamic universe (theory)



Alexander Friedman
(1888 -1925)

Alexander Friedman (1922)

Allow time-varying solutions for the cosmos

Two differential equations for R

$$\frac{3R'^2}{R^2} + \frac{3c^2}{R^2} - \lambda = \kappa c^2 \rho,$$

Evolving universe

Time-varying radius and density of matter

Distrusted by Einstein

$$\frac{R'^2}{R^2} + \frac{2RR''}{R^2} + \frac{c^2}{R^2} - \lambda = 0.$$

Georges Lemaître
(1894-1966)

Georges Lemaître (1927)

Theoretical universe of expanding radius

Agreement with emerging astronomical data

Also rejected by Einstein

$$\frac{1}{c^2} \left(\frac{dR}{dt} \right)^2 = \frac{A - R + \frac{\lambda}{3c^2} R^3}{R}$$



The watershed



Edwin Hubble (1889-1953)

Hubble's law (1929)

A redshift/distance relation for the galaxies

Linear relation: $h = 500 \text{ kms}^{-1}\text{Mpc}^{-1}$

Evidence of cosmic expansion?

RAS meeting (1930): Eddington, de Sitter

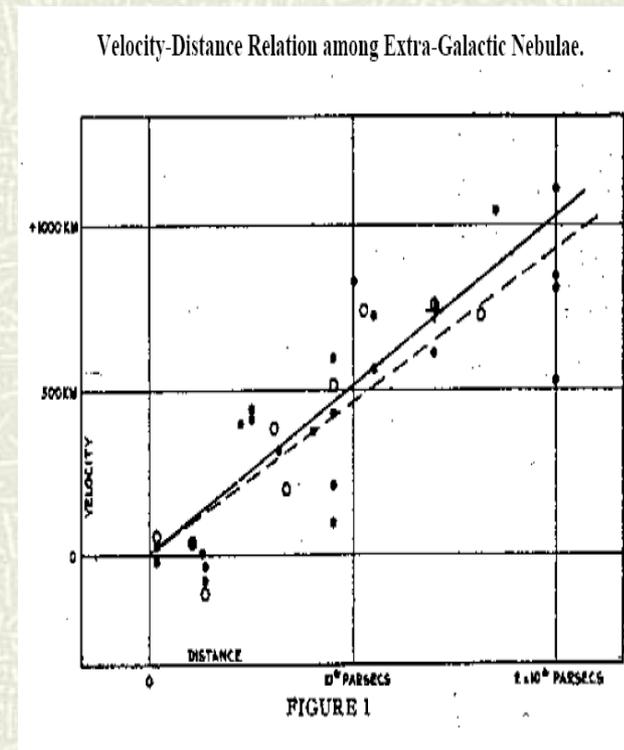
Friedman-Lemaître models circulated

Time-varying radius and density of matter

Einstein apprised

Sojourn at Cambridge (June 1930)

Sojourn at Caltech (Spring 1931)



The expanding universe (1930 -)

- **Eddington (1930, 31)**

*On the instability of the Einstein universe
Expansion caused by condensation?*

- **Tolman (1930, 31)**

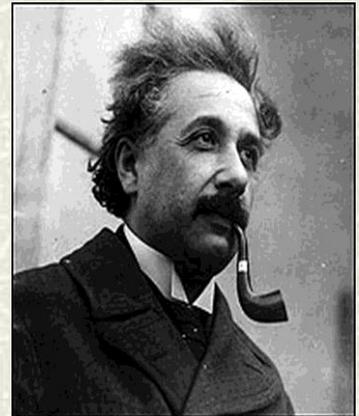
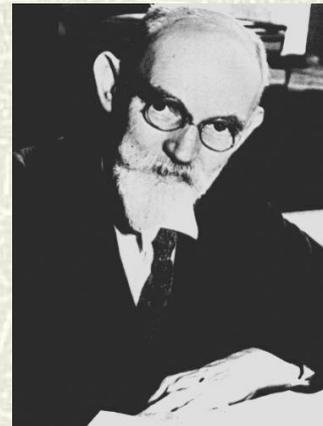
*On the behaviour of non-static models
Expansion caused by annihilation of matter ?*

- **de Sitter (1930, 31)**

*Further remarks on the expanding universe
Expanding universes of every flavour*

- **Einstein (1931, 32)**

*Friedman-Einstein model $k=1, \lambda=0$
Einstein-de Sitter model $k=0, \lambda=0$*



*Expanding models
No mention of origins*

Cosmic prediction IV: the big bang



- # **Lemaître 1931:** expanding U smaller in the past
- # Extrapolate to very early epochs
- # Extremely dense, extremely hot
- # Expanding and cooling ever since
- # 'Fireworks beginning' at $R = 0$?

Fr Georges Lemaître

Velocity-Distance Relation among Extra-Galactic Nebulae.

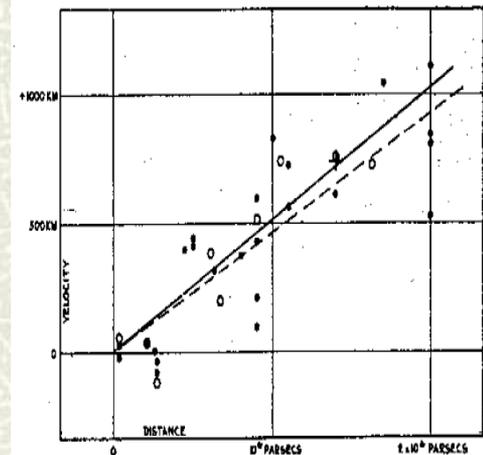


FIGURE 1

Not endorsed by community (1930-60)

Simplified models: timescale problem

Later called **'The big bang'**

The steady-state universe (1948)

‡ Expanding but unchanging universe

Perfect cosmological principle?

No extrapolation to early epochs

No beginning, no timescale paradox



Hoyle, Bondi, Gold (1948)

‡ Requires continuous creation of matter

Very little matter required

‡ Replace λ with creation term (Hoyle)

$$G_{\mu\nu} + C_{\mu\nu} = -k T_{\mu\nu}$$

‡ Improved version (1962)

$$G_{\mu\nu} + \lambda g_{\mu\nu} = k T (C_{\mu} + C_{\nu})$$



Hoyle and Narlikar (1962)

Steady-state vs big bang (1950-70)

Optical astronomy (1950s)

Revised distances to the nebulae (Baade, Sandage)

Timescale problem resolved

Radio-astronomy (1960s)

Galaxy distributions at different epochs

Cambridge 3C Survey (Ryle)

Nucleosynthesis of light elements

Alpher, Hermann and Gamow (1948)

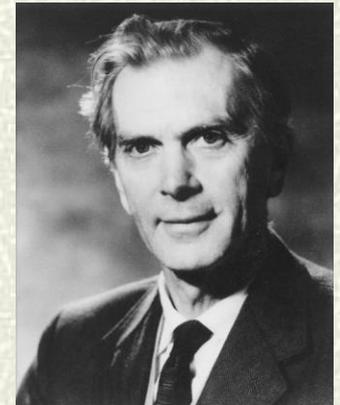
Cosmic microwave background (1965)

Microwave frequencies

Remnant of young, hot universe



Allen Sandage



Martin Ryle

Big bang puzzles

Characteristics of background radiation

Homogeneity, flatness, galaxy formation?(1970-80)

The theory of inflation (1981)

Exponential expansion within first second?

Which model of inflation?

Initial conditions?

Dark energy (1998)

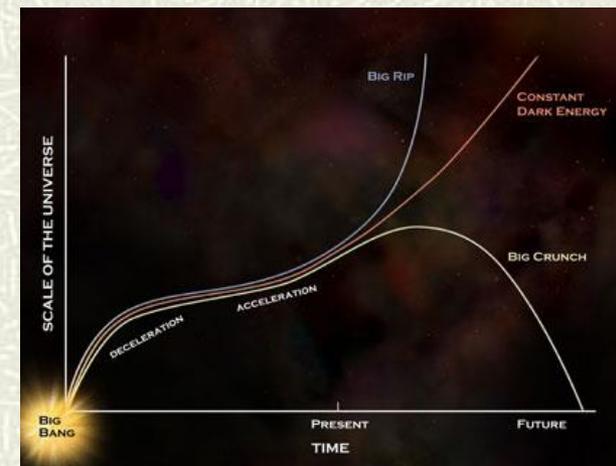
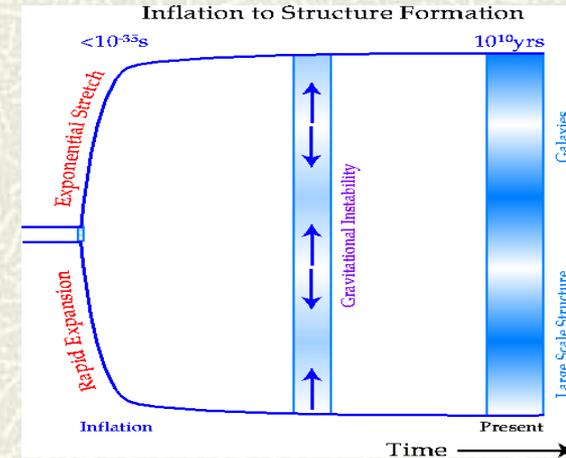
Observation of accelerated expansion

The return of the cosmological constant

Problems of interpretation

$$G_{\mu\nu} + \lambda g_{\mu\nu} = -\kappa T_{\mu\nu}$$

Nature of DE unknown



Relativity, astronomy and the universe: the first 100 years

Published 1915, 1916

A new theory of gravity

Classic predictions supported by observation

Perihelion of Mercury: bending of light by a star

Gravitational redshift

Cosmological predictions supported by observation

Black holes: gravitational waves

The expanding universe: the big bang

Relevant today

Modifications of theory?



Skeptical of extrapolations